

Optical control of spin-spin interaction in doped semiconductors

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Quantum Optoelectronics Theory
at MSU

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Collaborations

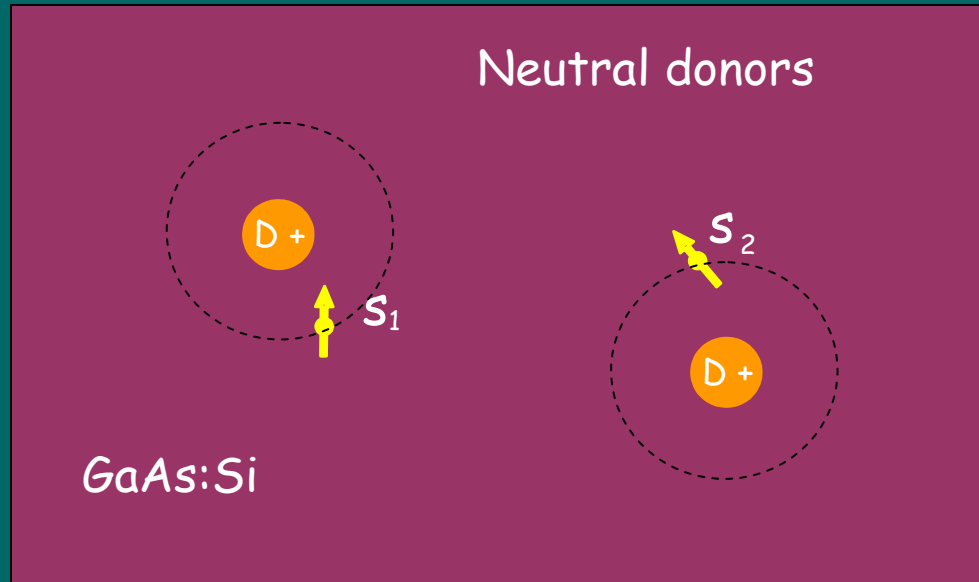
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UNIVERSITY

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Quantum Control of two spins

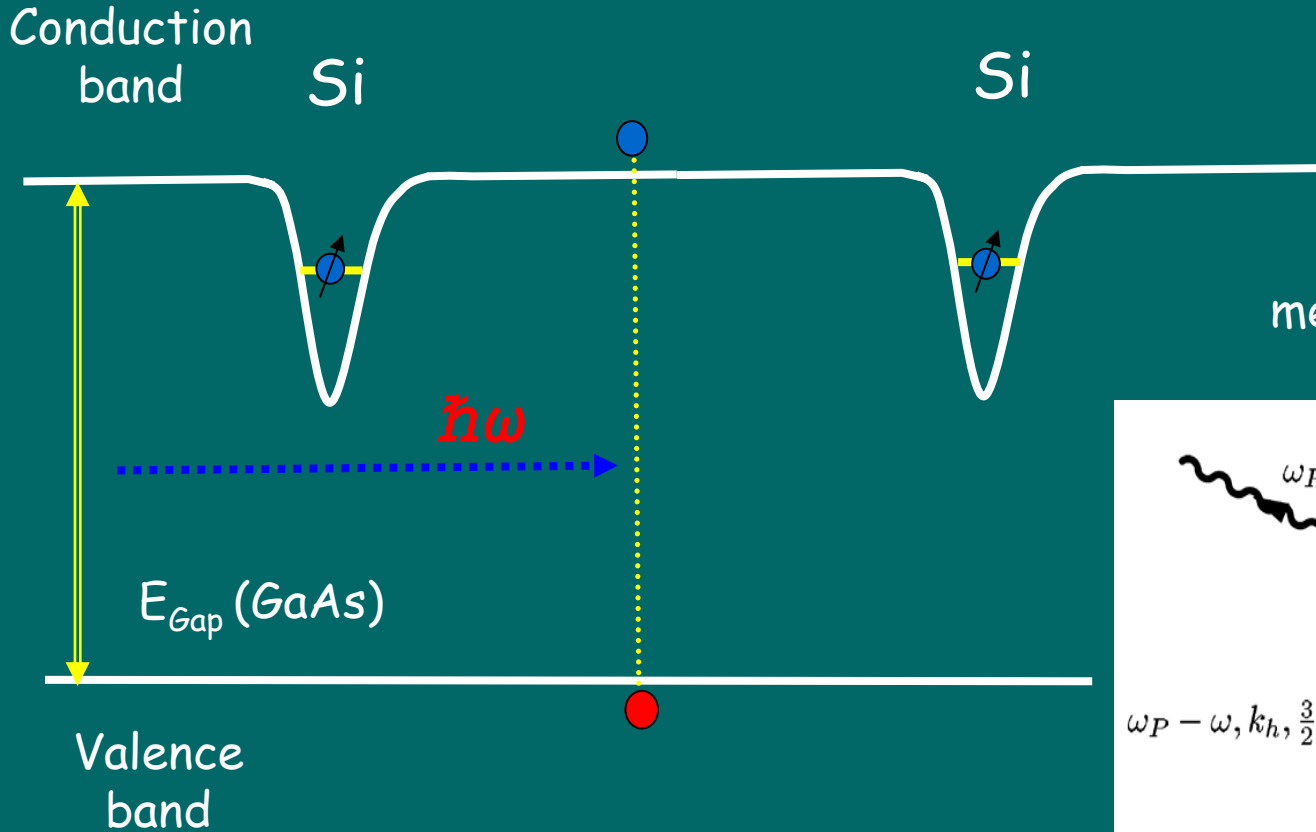


$$|\Psi\rangle = a(t)|\uparrow\uparrow\rangle + \beta(t)|\uparrow\downarrow\rangle + \gamma(t)|\downarrow\uparrow\rangle + \delta(t)|\downarrow\downarrow\rangle$$

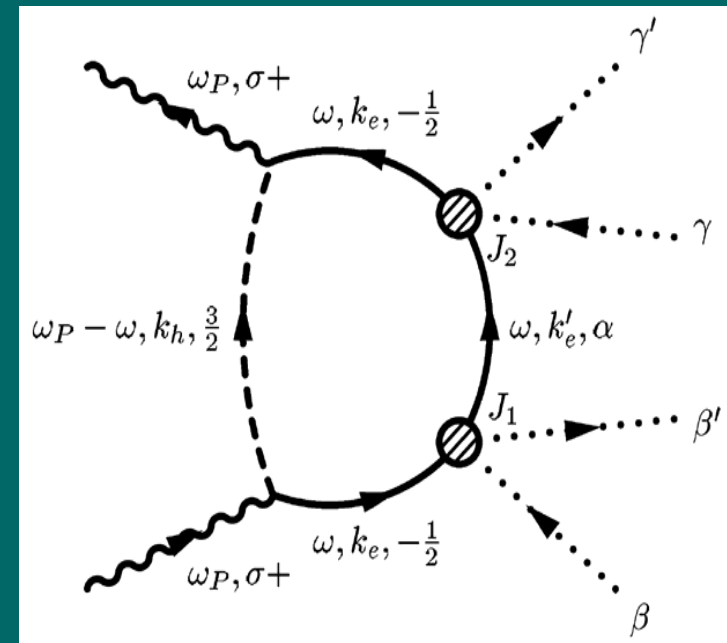
• Control Hamiltonian

$$H = H_0 + H_C(t, \sigma_1, \sigma_2, \dots)$$

Optical RKKY



Itinerant excitons mediate the interaction



C. Piermarocchi, P. Chen, L. J. Sham, G. D. Steel,
Phys. Rev. Lett. **89** 167402 (02)

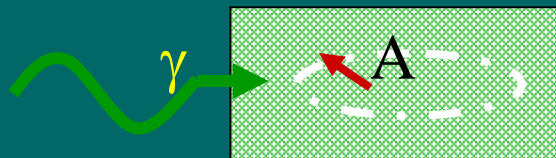
$$H = J_{\text{eff}} \mathbf{S}_1 \cdot \mathbf{S}_2$$

Beyond ORKKY

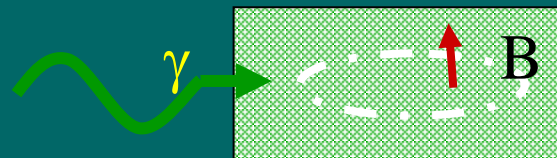
- Can we have Anti-ferromagnetic coupling?
- What is the effect of multiple scattering?
- What if the exciton is bound to the impurity?

Beyond second order in the exciton-spin coupling

G. F. Quinteiro and C. Piermarocchi, (to be published)



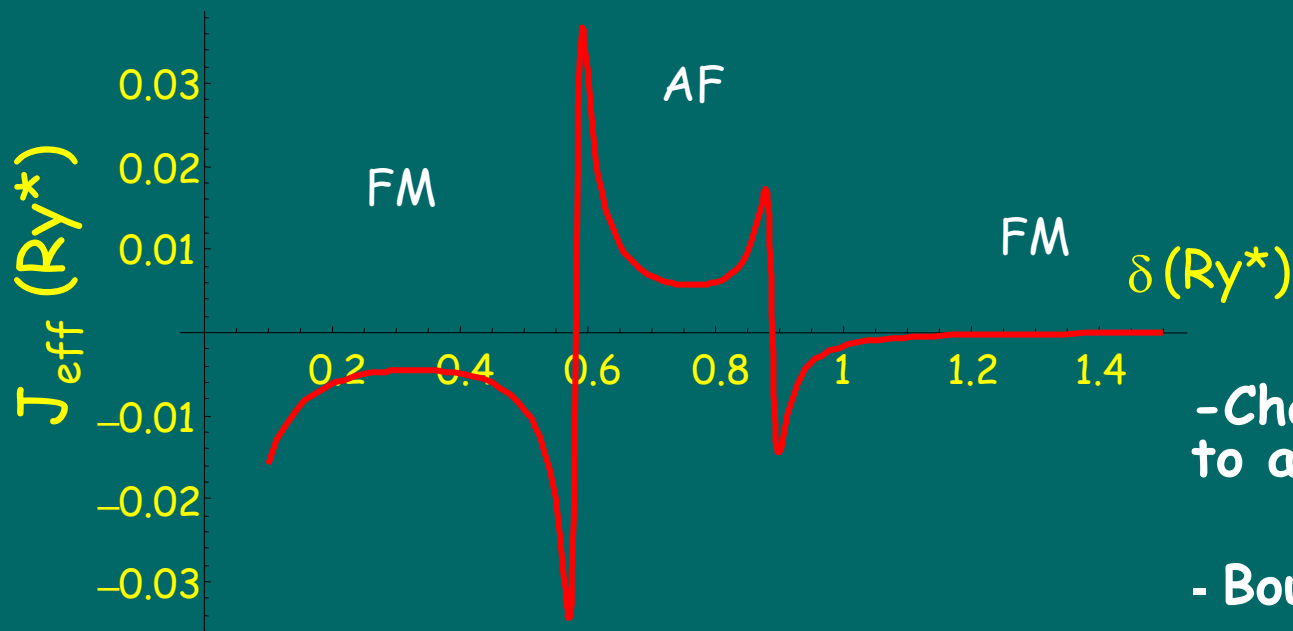
Solution for spin
A + exciton



Solution for spin
B + exciton



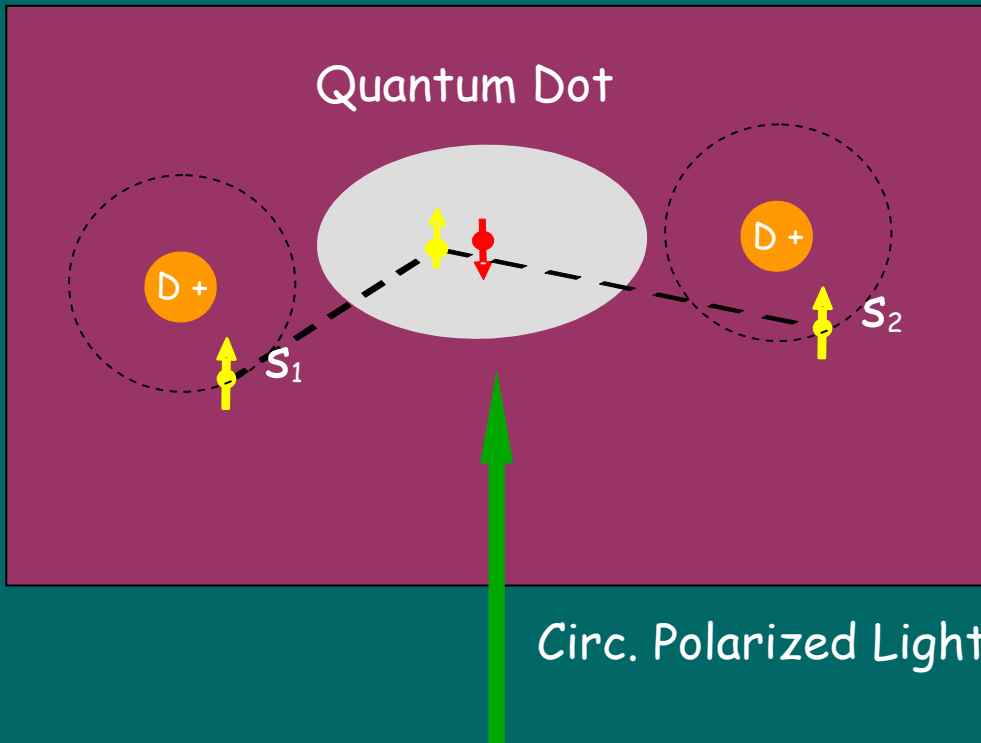
$$T = \frac{1}{1 - T^A G^0 T^B G^0} T^A [1 + G^0 T^B] + (A \rightleftharpoons B)$$



- Change from ferro
to antiferro

- Bound states.

Dynamics of quantum control



Initial State

$$|\uparrow\downarrow\rangle = \frac{|S\rangle + |T,0\rangle}{\sqrt{2}}$$

Follow the dynamics by numerical integration of the Liouville Equation

$$\frac{d\rho}{dt} = -i[H, \rho] + \mathcal{L}\{\rho\}$$

Dynamics of entanglement

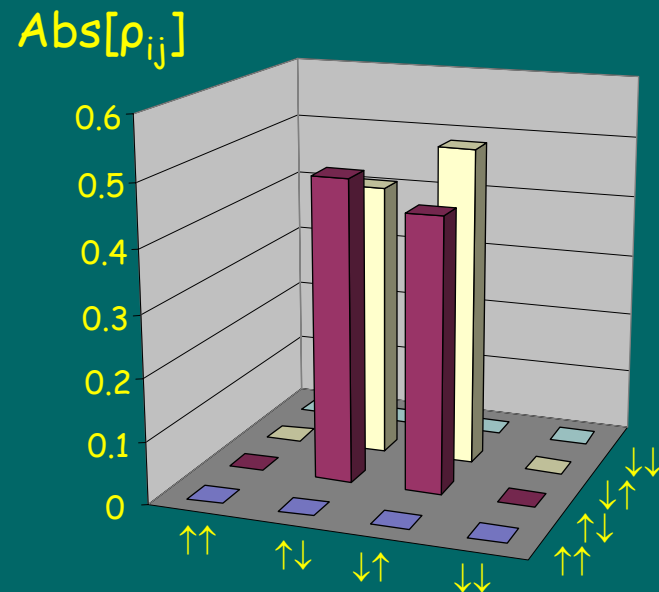
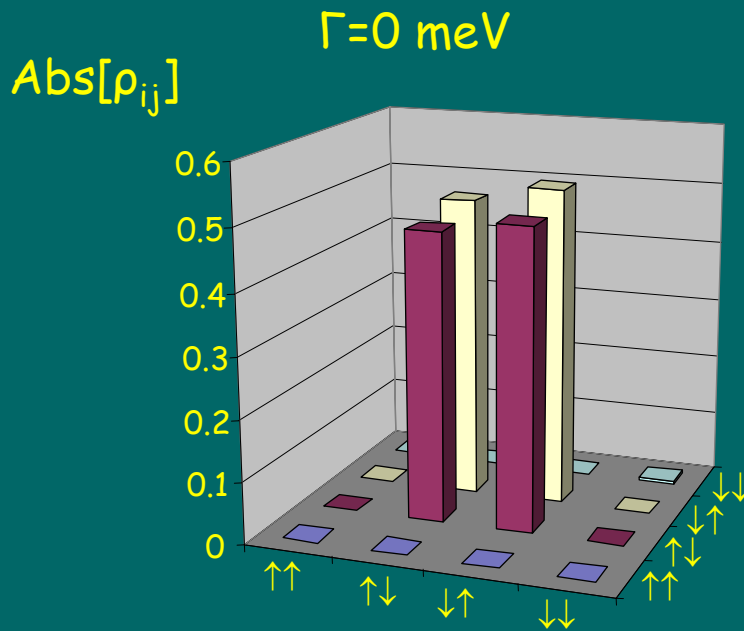
Effect of decoherence: radiative recombination

Decoherence

Off-resonant excitation

ρ after a Gaussian pulse

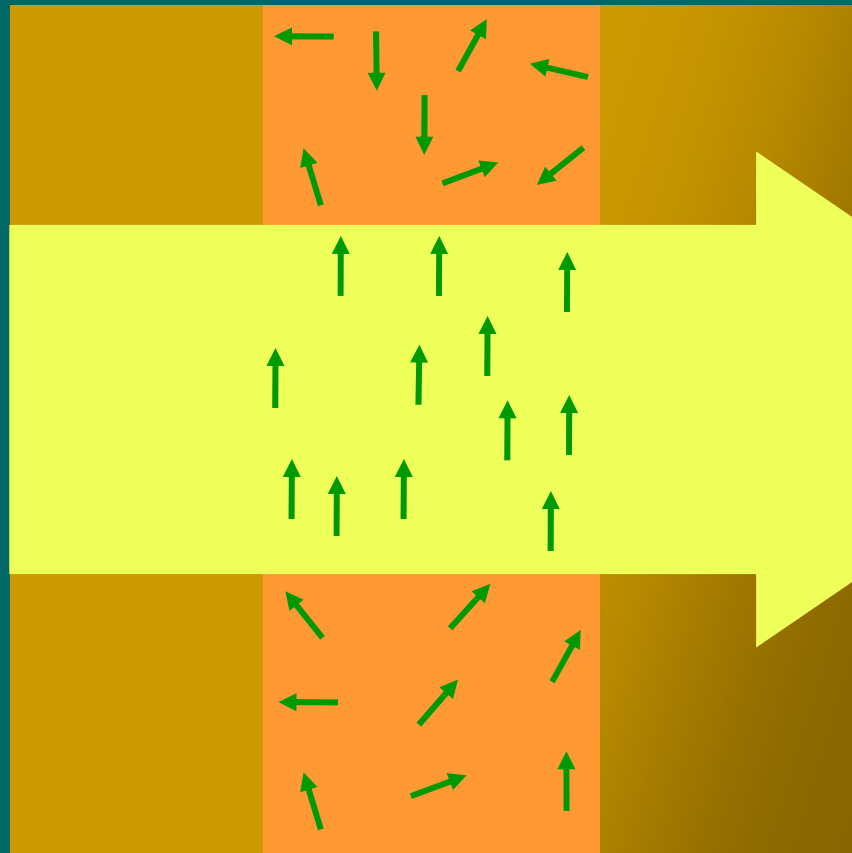
$\Gamma=0.3$ meV (Radiative Recombination)



$$|\Psi\rangle \sim \frac{|\uparrow\downarrow\rangle + i|\downarrow\uparrow\rangle}{\sqrt{2}}$$

Optical generation of entanglement !

Light-spin thermodynamics



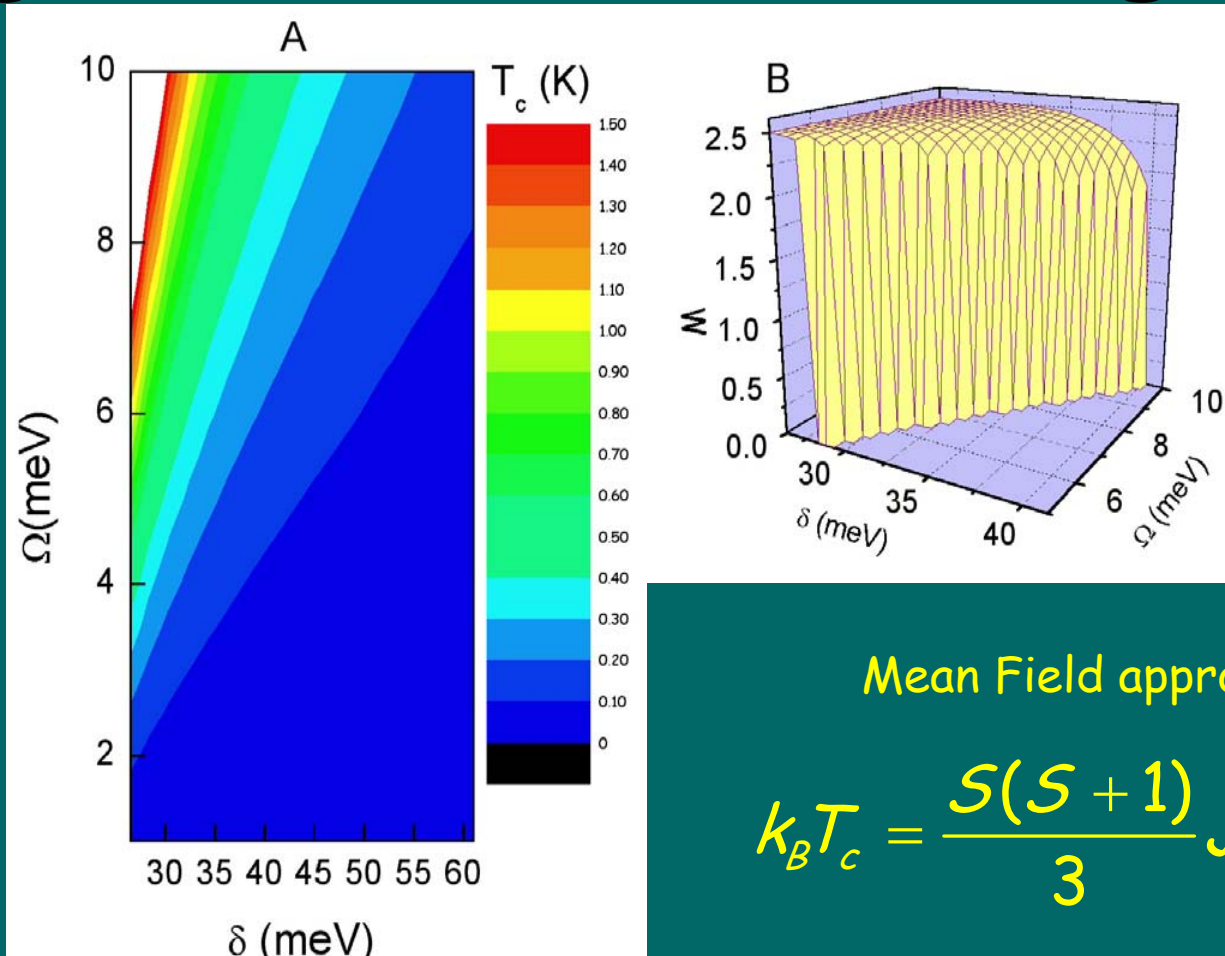
$$H = -J_{\text{ORKKY}}[\Omega] \sum_{ij} \mathbf{s}_i \cdot \mathbf{s}_j$$

Can we induce a PM/FM transition using coherent light?

ZnSe:Mn

J Fernandez-Rossier, C Piermarocchi, P Chen, LJ Sham, and AH MacDonald, (condmat/0312445, accepted in Phys. Rev. Lett.)

Light-induced ferromagnetism



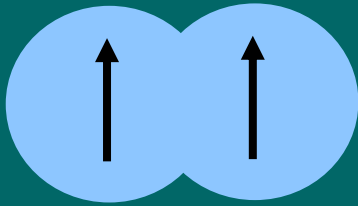
Mean Field approach

$$k_B T_c = \frac{S(S+1)}{3} J_{\text{ORKKY}} [\Omega]$$

Conclusions

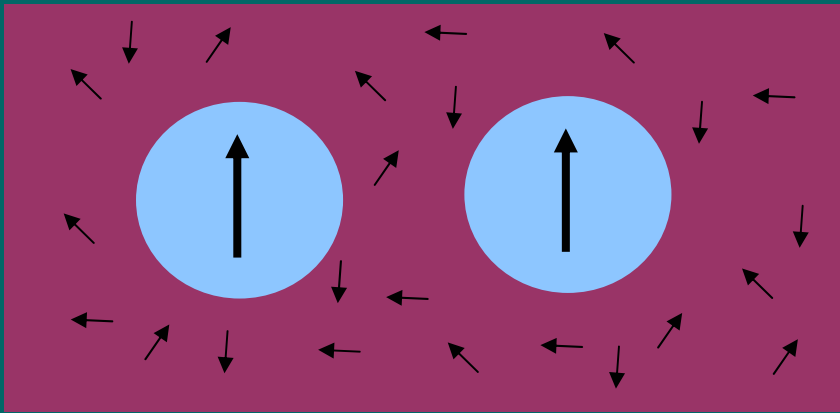
- Light can induce spin-spin interaction in a semiconductor.
- Strength, sign of the interaction are controllable.
- Dynamics and decoherence.
- Light-induced ferromagnetism.

Spin-spin interaction in Magnetism



DIRECT:

Charge distribution overlap
e. g. H_2



INDIRECT:

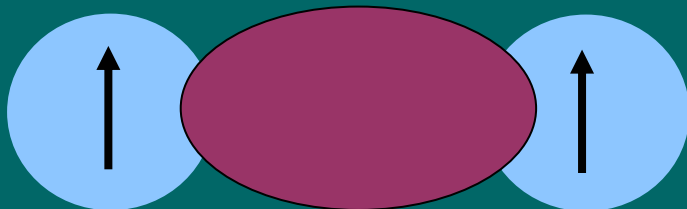
RKKY

Ex: Gd

or

Super-exchange

Ex: $CaMnO_3$



Dimensionality effects

$$H = -2J(R)\sigma_1 \cdot \sigma_2.$$

$$3\text{D} \quad J \sim \left(\frac{\Omega}{\Lambda}\right)^2 Ry^* \frac{e^{-2R/\kappa}}{R/\kappa},$$

$$2\text{D} \quad J \sim \left(\frac{\Omega^2}{\Lambda\delta}\right) Ry^* e^{-2R/\kappa},$$

$$1\text{D} \quad J \sim \left(\frac{\Omega}{\delta}\right)^2 Ry^* \left(1 + \frac{R}{\kappa}\right) e^{-2R/\kappa}.$$

Ω = Rabi energy

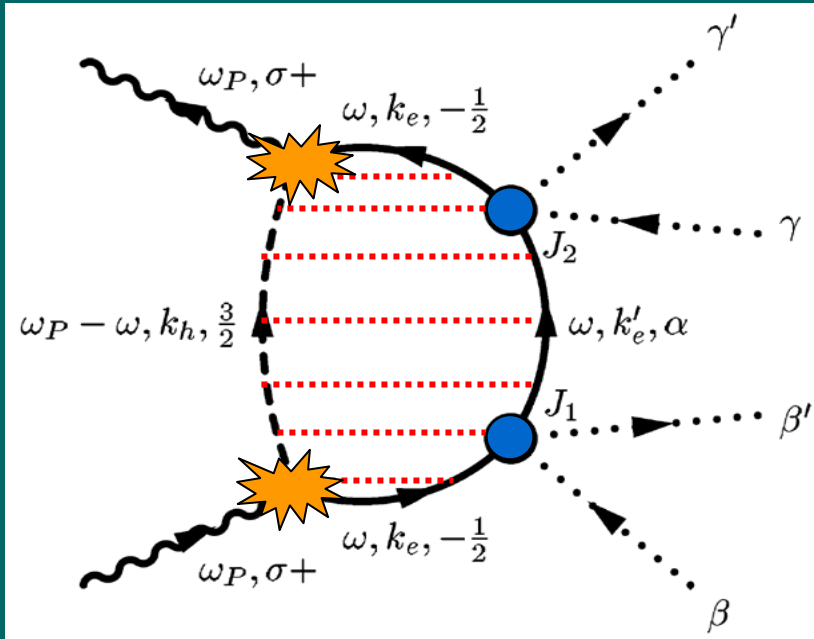
Λ = energy in dot

δ = detuning

R = interdot distance

$$\kappa = \sqrt{\frac{\hbar^2}{2m\delta}}$$

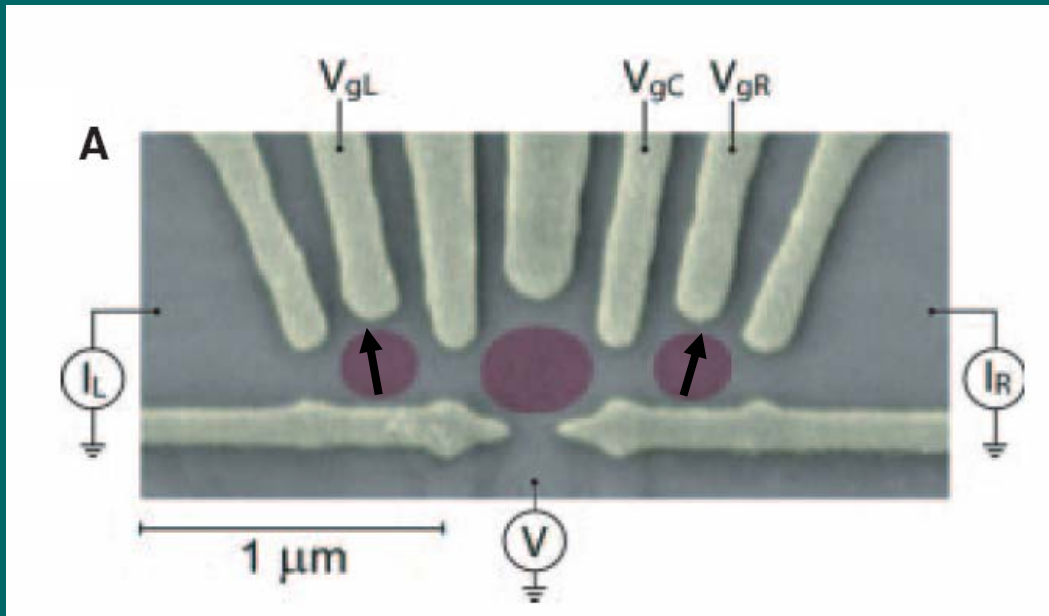
Excitonic effects



$$F_{1s,1s}(\mathbf{q}) = \int d\mathbf{r} e^{-i\frac{m_h}{M}\mathbf{q}\cdot\mathbf{r}} |\Psi_{1s}(r)|^2$$

$$J_{1s12}^d(R) = \frac{|\Omega(t)|^2}{16} j_1^d j_2^d \frac{1}{\Delta^3} |\Psi_{1s}(0)|^2 \int \frac{d^d \mathbf{q}}{(2\pi)^d} \frac{e^{-i\mathbf{q}\cdot\mathbf{R}}}{1 + (\lambda_M q)^2} |F_{1s,1s}(q)|^2$$

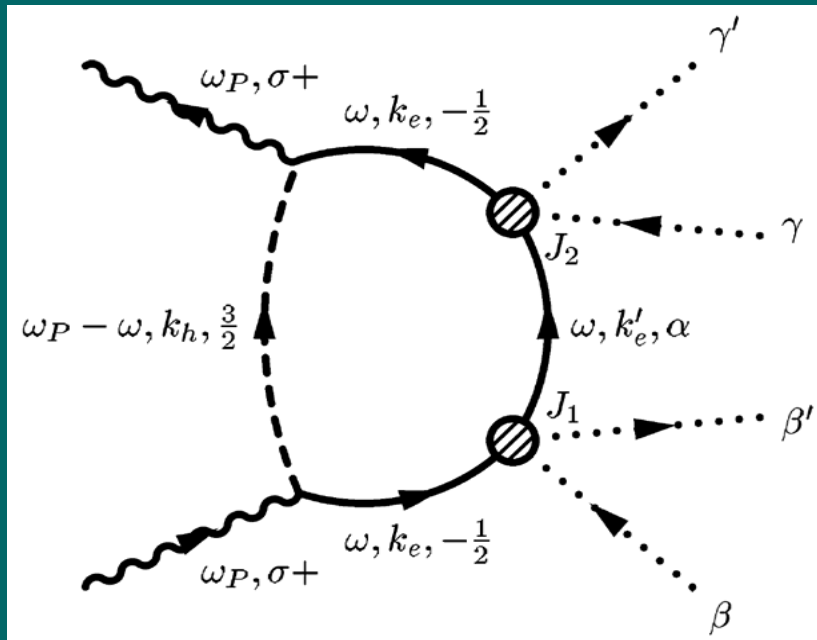
Controlled indirect spin-spin interaction in semiconductors



N.J. Craig et al.
Science April 23rd 2004

The central dot mediates an indirect RKKY-like spin-spin interaction between the lateral dots

Dynamic Spin-Stark shift



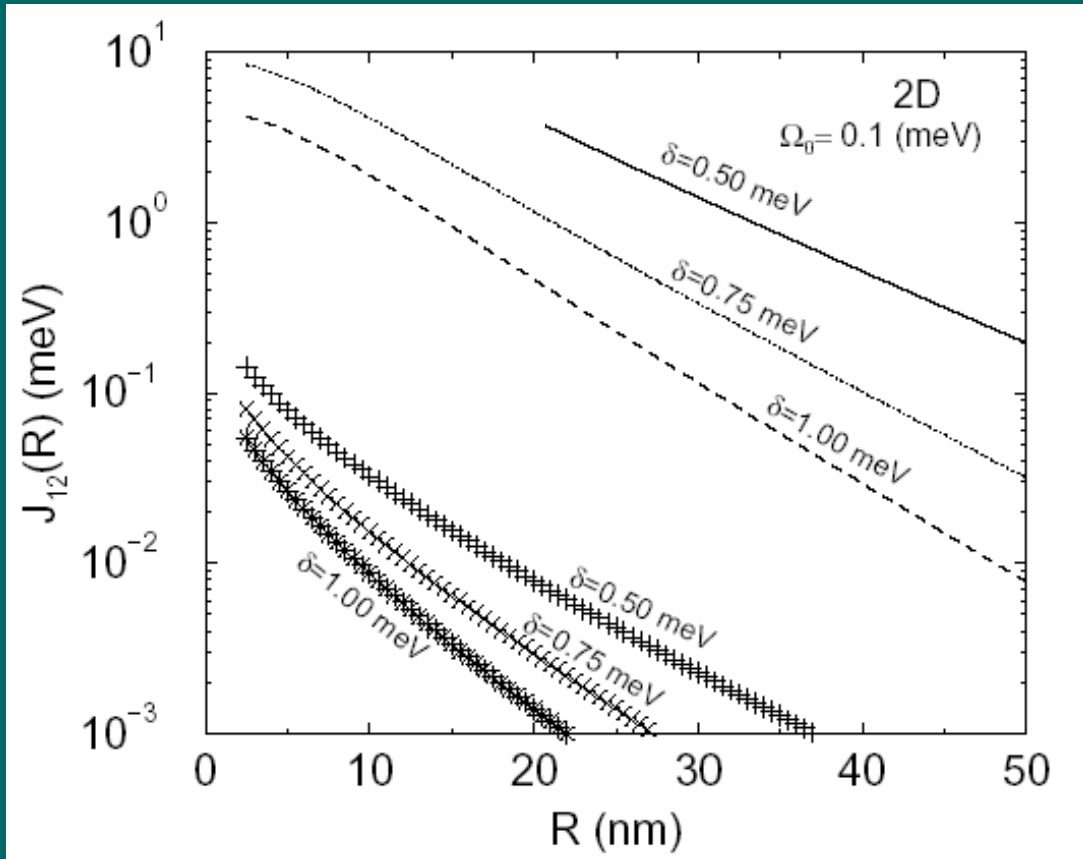
SPIN STRUCTURE:

$$-4J_{12}(\mathbf{S}^1 \cdot \mathbf{s})(\mathbf{S}^2 \cdot \mathbf{s}) - 4J_{12}(\mathbf{S}^2 \cdot \mathbf{s})(\mathbf{S}^1 \cdot \mathbf{s})$$

$$= -2J_{12}(\mathbf{S}^1 \cdot \mathbf{S}^2)$$

$$J_{12}(R) = \frac{|\Omega(t)|^2}{16} \iint \frac{d^d k}{(2\pi)^d} \frac{d^d k'}{(2\pi)^d} \frac{j_1^d j_2^d e^{-i(k-k')R}}{\left(\Delta + \frac{k^2}{2m_h} + \frac{k'^2}{2m_e} \right)^2 \left(\Delta + \frac{k^2}{2m_h} + \frac{k'^2}{2m_e} \right)}$$

Quantum Well



$$m_e^* = 0.07m,$$

$$m_h^* = 0.5m,$$

$$\xi = 300 \text{ \AA}$$

$$Ry^* = 10 \text{ meV}$$

$$a_B^* = 100 \text{ \AA}$$

Exponential decay of the interaction

$$K = \frac{\hbar}{\sqrt{2M_x\delta}}$$

Assuming $J_{kk'} = J v_k v_{k'}$



Exact analytical solution of the effective H of two localized spins

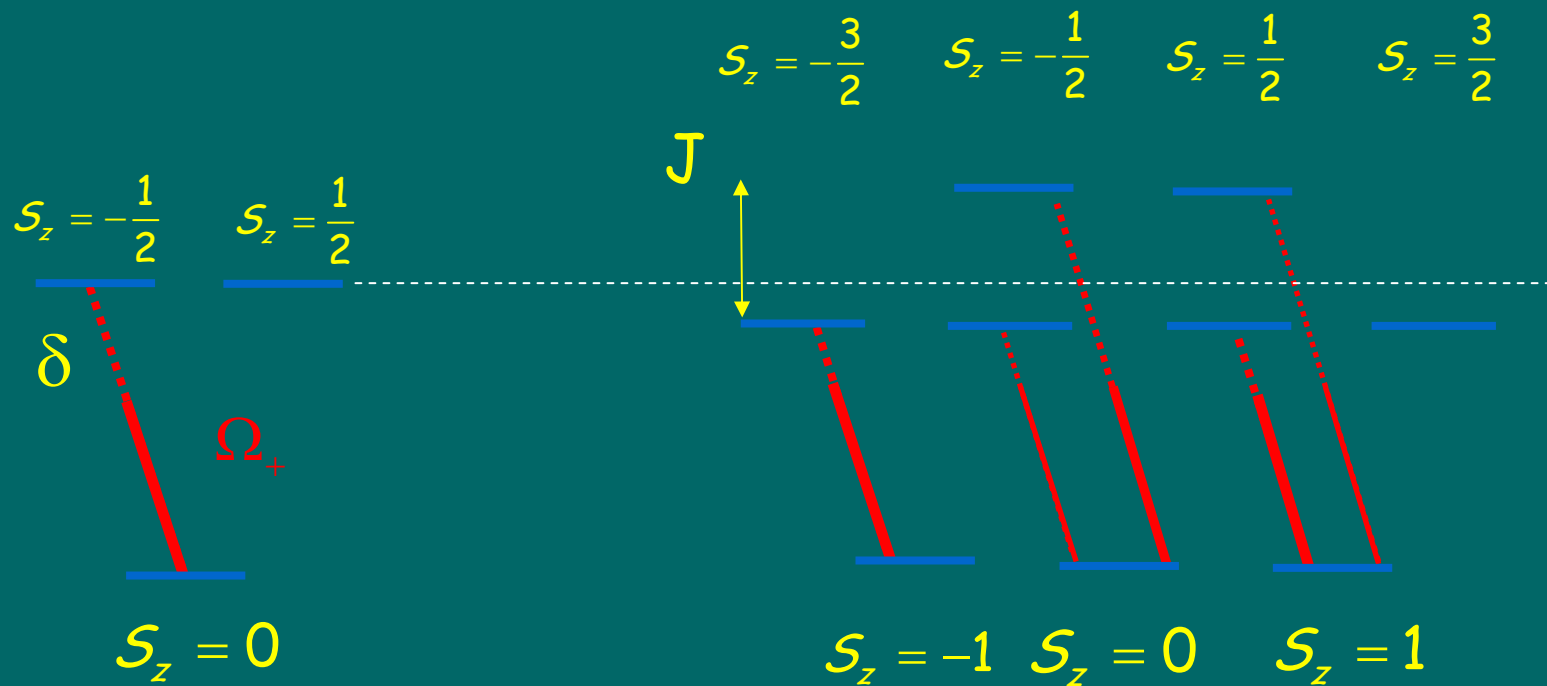
$$H_{\text{eff}} = -2 \left(\frac{\Omega}{\delta} \right)^2 |\psi_{1s}|^2 [A(s_1 + s_2) \cdot s + B s_1 \cdot s_2 + C(s_1 \times s_2) \cdot s]$$

A: effective magnetic field. Needs circular polarization

B: Heisenberg term

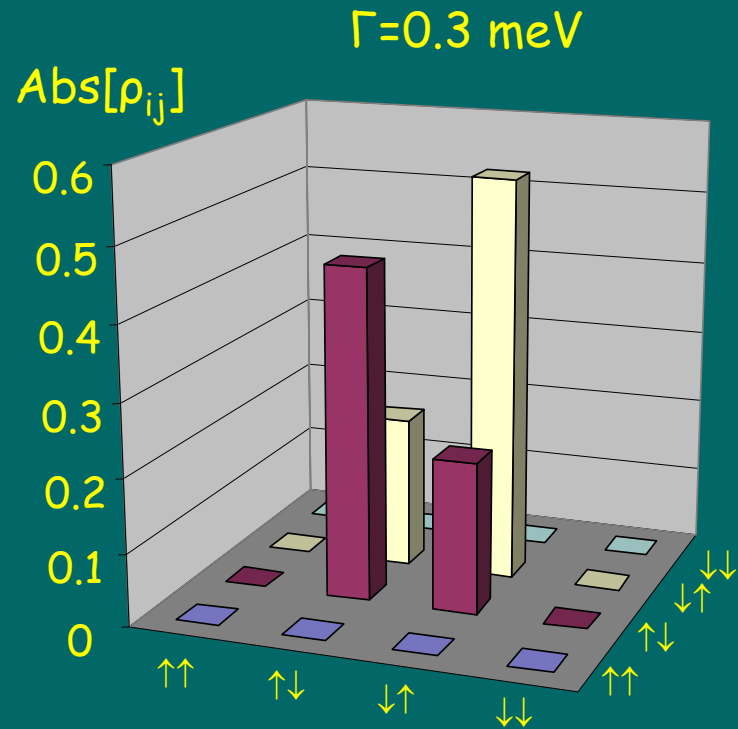
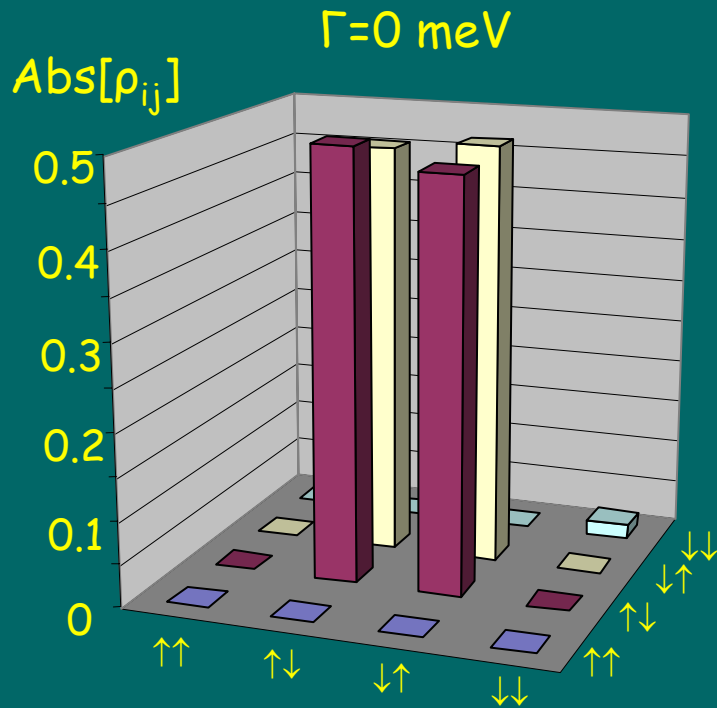
C: Dzyaloshinski-Moriya. Needs circular polarization + asymmetric coupling

Optical Spin-Stark Shift



Initial State $|\uparrow\downarrow\rangle = \frac{|S\rangle + |T,0\rangle}{\sqrt{2}}$

Radiative Recombination (b) resonance



$\delta=1$ meV

Resonant to the $J=3/2$ states

Nearly thermal in the $S_z=0$ subspace,
but still entangled in terms of Peres criterion