

Stochastic Growth in a Small World and Applications to Scalable Parallel Discrete-Event Simulations

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Abstract

We consider a simple stochastic growth model on a small-world network. The same process on a regular lattice exhibits kinetic roughening governed by the Kardar-Parisi-Zhang equation. In contrast, when the interaction topology is extended to include a *finite* number of random links for each site, the surface becomes macroscopically smooth. The correlation length of the surface fluctuations becomes finite and the surface grows in a mean-field fashion. Our finding provides a possible way to establish control *without* global intervention in non-frustrated agent-based systems. A recent application is the construction of a fully scalable algorithm for parallel discrete-event simulation.

Phase transitions in Small-World (SW) Networks

- Watts&Strogatz (1998): "... enhanced signal-propagation speed, computational power, and synchronizability".
- Finite number of random links per site (average degree is not extensive)
- Phase transition or phase ordering is possible even when random links are added to an originally *one-dimensional* substrate:
- Barrat&Weight (2000), Gitterman (2000), Kim et al. (2001), Herrero (2002), Jeong et al. (2003), Novotny and Wheeler (2004): Ising model on SW network
- Hong et al. (2002): XY-model and Kuramoto oscillators on SW network.
- Hastings (2003): general criterion for *mean-field-like* phase transitions for interacting systems on SW networks.

Synchronization in Parallel Discrete-Event Simulations

Parallelization for asynchronous dynamics

Paradoxical task:

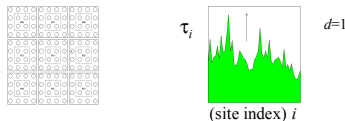
- (algorithmically) parallelize (physically) non-parallel dynamics

Difficulties:

- Discrete events (updates) are not synchronized by a global clock
- Traditional algorithms appear inherently serial (e.g., Glauber attempt one site/spin update at a time)

However, these algorithms are not inherently serial (Lubachevsky '87)

Two Approaches for Synchronization



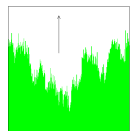
Optimistic (or speculative)

- PEs assume no causality violations
- Rollbacks to previous states once causality violation is found (extensive state saving or reverse simulation)
- Rollbacks can cascade ("avalanches")

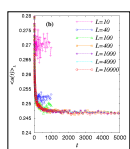
Conservative

- PE "idles" if causality is not guaranteed
- Utilization, $\langle u \rangle$: fraction of non-idling PEs

Basic Conservative Approach



- one-site-per PE, $N_{PE} = L^d$
- $t = 0, 1, 2, \dots$ parallel steps
- $\tau_i(t)$ local simulated time
- local time increments are iid exponential random variables
- advance only if $\tau_i \leq \min\{\tau_{im}\}$ (m : nearest neighbors)



Scalability modeling

- utilization (efficiency) $\langle u(t) \rangle$ (fraction of non-idling PEs)
- density of local minima
- width (spread) of time surface:

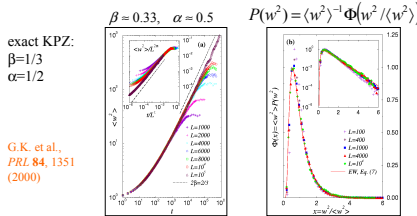
$$w^2(t) = \frac{1}{N_{PE}} \sum_{i=1}^{N_{PE}} [\tau_i(t) - \bar{\tau}(t)]^2$$

Simulating the Parallel Simulations

- Universality/roughness ($d=1$)

$$\langle w^2(t) \rangle_L \sim \begin{cases} t^{2\beta}, & \text{if } t \ll t_x \\ L^{2\alpha}, & \text{if } t \gg t_x \end{cases}, \quad t_x \sim L^z, \quad z = \alpha/\beta$$

Foltin et al., '94



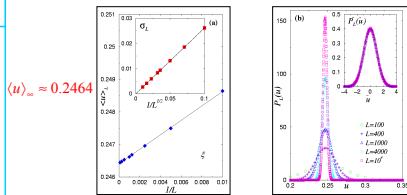
G.K. et al., PRL 84, 1351 (2000)

Utilization (Efficiency)

Finite-size effects for the density of local minima/average growth rate (steady state):

$$\langle u \rangle_L \equiv \langle u \rangle_\infty + \frac{const.}{L}$$

$$\sigma_L = \sqrt{\langle u^2 \rangle_L - \langle u \rangle_L^2} \sim 1/L^{1/2}$$



$\langle u \rangle_\infty = 0.2464$

Implications for Scalability

Simulation reaches steady state for $t \gg L^z$ (arbitrary d)

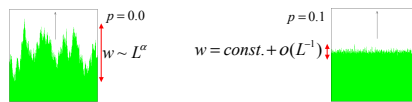
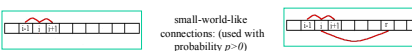
- Simulation phase: scalable $\langle u \rangle_L \equiv \langle u \rangle_\infty + \frac{const.}{L^{2(1-\alpha)}}$

$\langle u \rangle_\infty$ asymptotic average growth rate (simulation speed or utilization) is non-zero (Krug and Meakin, '90)

- Measurement (data management) phase: not scalable

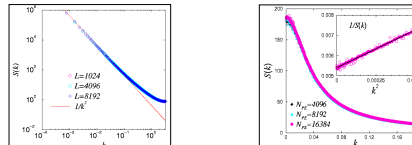
measurement at τ_{meas} (e.g., simple averages) rough (self-affine) synchronization landscape $w \sim L^\alpha$, $\xi \sim L$

Synchronization/Time-Horizon Control Via Small-World Communication Network Design



1d (ring) steady-state structure factor: (Fourier transform) $S(k) \propto \langle \tau_k \tau_{-k} \rangle = \langle |\tau_k|^2 \rangle$

$$S(k) \sim \frac{1}{k^2}$$

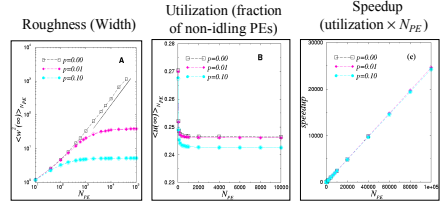


Utilization Trade-off/Scalable Data Management

$$\partial_t \tau = -\Sigma(p)\tau + \frac{\partial^2 \tau}{\partial x^2} + \dots + noise$$

G.K. et al., Science 299, 677 (2003)

effective relaxation to the mean facilitated by the SW links



Edwards-Wilkinson Model on a Small-World Network

$$\partial_t h_i = -(2h_i - h_{i+1} - h_{i-1}) - \sum_{j=1}^N J_{ij}(h_i - h_j) + \eta_i(t)$$

$$\partial_t h_i = -\sum_{j=1}^N \Gamma_{ij} h_j + \eta_i(t) \quad \Gamma_{ij}^o = 2\delta_{ij} - \delta_{i,j+1} - \delta_{i,j-1}$$

$$\Gamma = \Gamma^o + V$$

"soft" SW network (Erdős-Rényi network on top of the ring):



$$J_{ij} = \begin{cases} 1 & \text{with probability } p/N \\ 0 & \text{with probability } 1-p/N \end{cases}$$

"hard" SW network:



$J_{ij} = 0$ or p : $N/2$ random links are selected, such that each site has exactly one random link of strength p (in addition to n.n.)

Width from exact numerical diagonalization:

$$\{\lambda_j\}_{j=0}^{N-1} \text{ eigenvalues of } \Gamma_{ij} \quad (\lambda_0 = 0)$$

$$\langle w^2 \rangle_N = \frac{1}{N} \sum_{i=1}^N \langle (h_i - \bar{h})^2 \rangle = \frac{1}{N} \sum_{i=1}^{N-1} \frac{1}{\lambda_i}$$

for a single realization of the random network

$$[\langle w^2 \rangle_N] \text{ averaged over network realizations ("disorder-averaged" width)}$$

Impurity-averaged perturbation theory

Kozma et al. PRL 92, 108701 (2004).

$$[G]^{-1} = \Gamma^o + \Sigma \quad G = \Gamma^{-1} \quad G^o = \Gamma^{o-1}$$

$$[\langle w^2 \rangle_N] = [\langle \Gamma^{-1} \rangle_{ii}] \quad \text{self-energy (effective interaction due to random links)}$$

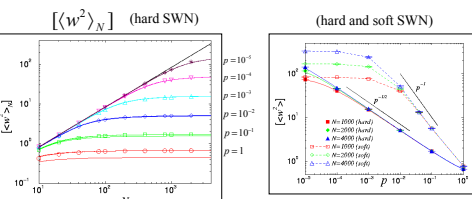
"soft" network: $\Sigma \sim p^2 + \dots$ [see also Monasson, EPJB 12, 555 (1999) in the context of diffusion on SWN]

"hard" network: $\Sigma \sim p - \frac{1}{2} p^{3/2} + \dots$

$$\xi \sim \frac{1}{\sqrt{\Sigma}} \quad [\langle w^2 \rangle_N] \sim \frac{1}{2\sqrt{\Sigma}}$$

finite correlation length for any $p > 0$ finite width

Comparison of exact numerical diagonalization of Γ_{ij} with the results of the impurity-averaged perturbation theory



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Summary

- Synchronizability of large-scale non-frustrated agent-based systems with SW network: application to construct fully scalable parallel simulations *without* global synchronizations
- Spectrum of the coupling matrix exhibits a gap/pseudo-gap, yielding a *finite width* for stochastic growth on a small-world network for all $p > 0$